

2011 Day 2 Question 3A (GPP)

a. $\frac{\partial n}{\partial t} + \nabla \cdot (n v) = 0$

$\dot{n} + n_0 v' = 0$

Linearize, assume 1D

$\nabla^2 \phi = -4\pi g n$

$\phi'' = -4\pi g n$

$m \frac{dv}{dt} = -g \phi'$

$m \dot{v} = -g \phi'$

$\Rightarrow m v' = -g \phi''$

$-\frac{m}{n_0} \ddot{n} = 4\pi g^2 n \Rightarrow \left(\frac{4\pi g^2 n_0}{m} + \partial_t^2 \right) n = 0$

Fourier Transform in time: $(\omega_p^2 - \omega^2) n = 0 \quad \omega^2 = \omega_p^2$

b. $E \phi' = -4\pi g n$, $E = E_0 \cos(kx - \omega t)$

$-k E_0 \sin(kx - \omega t) = -4\pi g n \rightarrow n(x, t=0) = \frac{k E_0 \sin(kx)}{4\pi g}$

$m \dot{v} = -g E_0$

$\dot{v} = -\frac{g}{m} E_0 \cos(kx - \omega t)$

$v = -\frac{g \omega}{m} E_0 \sin(kx - \omega t) \Big|_{t=0} = -\frac{g \omega}{m} E_0 \sin(kx)$

c. If $E(t \rightarrow \infty) = 0$ is nonzero, it has completed the damping of f .



$N_{\text{trans}} \approx \Delta v \int_0^{\frac{\omega}{k} - \Delta v} f_0 \rightarrow N_{\text{trans}} \approx \Delta v \int_0^{\frac{\omega}{k} + \Delta v} f_0$

$f_0(\frac{\omega}{k} \pm \Delta v) \approx f_0(\frac{\omega}{k}) \pm \Delta v f_0'(\frac{\omega}{k})$

$N = N_a - N_b \approx \Delta v [f_0(\frac{\omega}{k} - \Delta v) - f_0(\frac{\omega}{k} + \Delta v)] = 2 \Delta v^2 |f_0'(\frac{\omega}{k})|$

Energy gained $\approx \frac{1}{2} m (\frac{\omega}{k})^2 - (\frac{\omega}{k} - \Delta v)^2 \approx m (\frac{\omega}{k}) \Delta v$

Total energy transferred $\approx N \cdot E_{\text{quanta}} = 2 m (\frac{\omega}{k}) |f_0'(\frac{\omega}{k})| \Delta v^3$

where $\Delta v = v_{gr} = \sqrt{2 E_0 / m k}$

To not extinguish wave, must have $\frac{E^2}{8\pi} > 2 m (\frac{\omega}{k}) |f_0'(\frac{\omega}{k})| \Delta v^3$

[Missing factor of n_0 on RHS. If had, could get $\omega_p^2 \omega \approx \omega_p^3$ and have result.]

(why?)

2011 Day 2 Q3A (GPP)

a. $\frac{\partial n}{\partial t} + \nabla \cdot (n\mathbf{v}) = 0 \rightarrow \frac{\partial \tilde{n}}{\partial t} + n_0 \nabla \cdot \tilde{\mathbf{v}} = 0$

$m n \frac{d\mathbf{v}}{dt} = -q \nabla \phi$

$m \frac{\partial \tilde{\mathbf{v}}}{\partial t} = -q \nabla \tilde{\phi}$

$\nabla^2 \phi = -4\pi q n$

$\nabla^2 \tilde{\phi} = -4\pi q \tilde{n}$

$m \frac{\partial}{\partial t} (\nabla \cdot \tilde{\mathbf{v}}) = -q \nabla^2 \tilde{\phi} = 4\pi q^2 \tilde{n}$

$m \frac{\partial}{\partial t} \left(-\frac{1}{n_0} \frac{\partial \tilde{n}}{\partial t} \right) = 4\pi q^2 \tilde{n}$

$(\partial_t^2 + \omega_p^2) \tilde{n} = 0$

$\omega_p^2 = \frac{4\pi n e^2}{m}$

$(-\omega^2 + \omega_p^2) = 0 \rightarrow \omega^2 = \omega_p^2$

b. $\tilde{\mathbf{E}} = E_0 \cos(kx - \omega t)$, ~~$\tilde{\phi} = -\frac{\partial \tilde{\phi}}{\partial x} = k E_0 \sin(kx - \omega t)$~~

$\nabla^2 \tilde{\phi} = -4\pi q \tilde{n}$

$-\nabla \cdot \tilde{\mathbf{E}} = -4\pi q \tilde{n}$

$k E_0 \sin(kx - \omega t) = -4\pi q \tilde{n}(x, t)$

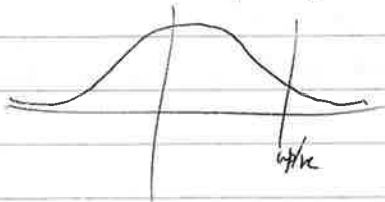
$\tilde{n}(x, t=0) = -\frac{k E_0}{4\pi q} \sin(kx)$

$m \frac{\partial \tilde{\mathbf{v}}}{\partial t} = -q \nabla \tilde{\phi} = q \tilde{\mathbf{E}} \Rightarrow \frac{\partial \tilde{\mathbf{v}}}{\partial t} = \frac{q}{m} E_0 \cos(kx - \omega t)$

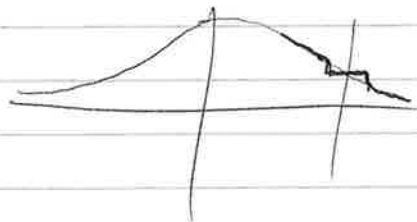
$\tilde{\mathbf{v}}(x, t) = -\frac{q}{m \omega p} E_0 \sin(kx - \omega t)$, $\tilde{\mathbf{v}}(x, t=0) = \frac{-q E_0}{m \omega p} \sin(kx)$

c. If wave not completely extinguished, the distribution will be fully flattened near resonance.

Also damps at $\omega = \omega_p$



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suppose x already integrated out
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 $v_p \approx \frac{\omega_p}{k}$

d. $N_{\text{trans}} \approx \Delta v \int_0^L (v_p - \Delta v)$

$N_{\text{refl}} \approx \Delta v \int_0^L (v_p + \Delta v)$

$N = N_{\text{trans}} - N_{\text{refl}} = \Delta v \int_0^L (v_p - \Delta v) - \int_0^L (v_p + \Delta v)$

Assume resonance near tail of distribution, f slowly varying

$f_0(v_p \pm \Delta v) \approx f_0(v_p) \pm \Delta v f_0'(v_p) = f_0(v_p) \pm \Delta v |f_0'(v_p)|$

$N = 2(\Delta v)^2 |f_0'(v_p)|$

Change in energy for single particle accelerated from $v_p - \Delta v$ to v_p :

$\Delta E = \frac{1}{2} m_e (v_p^2 - (v_p - \Delta v)^2) = m_e v_p \Delta v$

Work done by the wave $W = 2 \int \vec{E} \cdot d\vec{x} \approx \frac{E_0 e}{k}$

Assume $W \approx \Delta E \rightarrow \Delta v = \frac{E_0 e}{m_e k v_p} = \frac{e E_0}{m_e \omega_p} \equiv V_{\text{trapped}}$

To not extinguish wave, must have ^{total} wave energy $>$ total work done:

$\frac{E_0^2}{4\pi} > N \Delta E = 2 m_e v_p (\Delta v)^2 |f_0'| = 2 m_e \frac{\omega_p}{k}$

In frame of phase velocity, $m \ddot{x} = -e E_0 \cos(kx)$

$\rightarrow \ddot{x} + \frac{e E_0 k}{m} x = 0 \rightarrow \omega_B = \sqrt{\frac{e E_0 k}{m}}, v_{tr} = \frac{\omega_B}{k} = \sqrt{\frac{e E_0}{m k}}$

And assume $\Delta v \approx v_{tr}$.

$\frac{E_0^2}{4\pi} > N \Delta E = 2 m_e \frac{\omega_p}{k} v_{tr}^2 \left| \frac{\partial f}{\partial v} \left(\frac{\omega_p}{k} \right) \right|$

$\frac{1}{4\pi} \left(\frac{m k v_{tr}^2}{e} \right)^2 > 2 m \frac{\omega_p}{k} v_{tr}^2 \left| \frac{\partial f}{\partial v} \right|$

$v_{tr} > \sqrt{\frac{2\pi e^2}{m^2}} \frac{\omega_p}{k^2} \left| \frac{\partial f}{\partial v} \right|$ almost right...

$\rightarrow \omega_p^2$ if there was factor of density...