

2012 Day 1 Question 2 (Waves)

$$\frac{\partial f}{\partial t} + \mathbf{v} \cdot \nabla f + \mathbf{F} \cdot \nabla_{\mathbf{p}} f = 0$$

$$\frac{\partial f_1}{\partial t} + \mathbf{v} \cdot \nabla f_1 + \tilde{\mathbf{F}} \cdot \nabla_{\mathbf{p}} f_0 = 0 \quad \tilde{\mathbf{F}} = g \mathbf{E} = -\epsilon \nabla \phi$$

$$(-i\omega + ik \cdot \mathbf{v}) f_1 = + g i \hbar \epsilon k \cdot \nabla_{\mathbf{p}} f_0$$

$$f_1 = \frac{g \epsilon k \cdot \nabla_{\mathbf{p}} f_0}{k \cdot \mathbf{v} - \omega}$$

$$\nabla \cdot \tilde{\mathbf{E}} = -\nabla^2 \phi = -(ik)^2 \phi = \frac{4\pi g}{k^2} = \frac{4\pi}{k^2} \int f_1 d^3 p$$

$$\rightarrow \phi = \frac{4\pi \epsilon}{k^2} \int f_1 d^3 p = \frac{4\pi \epsilon^2}{k^2} \int \frac{\phi}{k \cdot \mathbf{v} - \omega}$$

$$\epsilon = 0 = 1 - \frac{4\pi \epsilon^2}{k^2} \int \frac{k \cdot \nabla_{\mathbf{p}} f_0}{k \cdot \mathbf{v} - \omega} d^3 p$$

$$\nabla_{\mathbf{p}} f_0 = \frac{c}{T} f_0 \rightarrow \epsilon = 1 + \frac{4\pi \epsilon^2 c}{k^2 T} \int \frac{k \cos \theta f_0}{k \cos \theta v - \omega} p^2 \sin \theta d\theta dp (2\pi)$$

Ultrarelativistic,  $v \approx c$  (although we integrate over all  $p(v)$ ),

$$\epsilon \approx 1 + \frac{4\pi \epsilon^2}{k^2 T} \int_0^\infty f_0(p) 2\pi p^2 dp \int_0^\pi \frac{\cos \theta}{\cos \theta - v} d\theta \sin \theta \left( v = \frac{\omega}{kc} \right)$$

$$\begin{aligned} b. \epsilon &= 1 + \frac{2\pi \epsilon^2 \lambda_0}{k^2 T} \int_0^\pi \frac{\cos \theta \sin \theta d\theta}{\cos \theta - v} \quad \text{let } s = \sin \theta \\ &= 1 + \frac{1}{2k^2 \lambda_0^2} \int_{-1}^1 \frac{s}{s - v} ds \end{aligned}$$

For  $v \in (-1, 1)$  the denominator goes to zero, so the integral diverges.

We are then compelled to interpret the integral with the Landau contour, namely bending the integration contour below the singularity at  $\ln \omega = 0$ .

Technically our Laplace Transform above was for  $\ln \omega > 0$ , so to extend to  $\ln \omega \leq 0$  as well, we need to use the Landau Contour.

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Note  $|U|$  since if  $U > 0$ ,

$$\begin{aligned} b. \int_{-1}^1 \frac{s}{s+U} ds &= \int_{-1}^1 \frac{s-U+U}{s-U} ds = \int_{-1}^1 1 + \frac{U}{s-U} ds \\ &= 2 + U \int_{-1}^1 \frac{ds}{s-U} \end{aligned}$$

For  $|U| > 1$ , not divergent  $\int$  can be evaluated as usual.  
 $= 2 + U \log \left( \left| \frac{1-U}{1+U} \right| \right)$

For  $|U| < 1$ , must take residue.  $1^{st}$  order pole:  $i\pi/U$   
 $= 2 + U \log \left| \frac{1-U}{1+U} \right| + i\pi/U$

$$\rightarrow \operatorname{Re} E = 1 + \frac{1}{k^2 \lambda_0^2} + \frac{1}{2k^2 \lambda_0^2} U \log \left| \frac{1-U}{1+U} \right|$$

$$\operatorname{Im} E = \frac{i\pi/U}{2k^2 \lambda_0^2}$$

Note  $|U| \rightarrow 1$ ,  $\operatorname{Re} E$  has logarithmic singularity.

c. Some approximations first:

$$U \ll 1 \quad \begin{cases} \log(1-U) \approx -U - \frac{U^2}{2} = -U \left( 1 + \frac{U}{2} \right) \\ \log(1+U) \approx U - \frac{U^2}{2} = U \left( 1 - \frac{U}{2} \right) \\ \rightarrow U \log \left( \frac{1-U}{1+U} \right) \approx -U^2 \end{cases}$$

$$U \gg 1 \quad \log \left| \frac{1-U}{1+U} \right| \stackrel{\text{w/o log}}{\approx} \log \left| \left( \frac{1-U}{U} - 1 \right) \right| \approx \log \left| (1-\frac{1}{U})(1-\frac{1}{U}) \right| \approx \log \left| 1 - \frac{2}{U} \right|$$

$$\log \left( 1 - \frac{2}{U} \right) \approx -\frac{2}{U} - \frac{1}{2} \left( \frac{2}{U} \right)^2 = -\frac{2}{U} \left( 1 + \frac{2}{U} \right) \quad \rightarrow U \log \left| \frac{1-U}{1+U} \right| \approx -2 \left( 1 + \frac{2}{U} \right)$$

so for  $U \ll 1$ ,  $\operatorname{Re} E = 1 + \frac{1}{k^2 \lambda_0^2} - \frac{U^2}{k^2 \lambda_0^2} \approx 1 + \frac{1}{k^2 \lambda_0^2}$  at  $U=0$ .

$U \gg 1$ ,  $\operatorname{Re} E = 1 + \frac{1}{k^2 \lambda_0^2} - \frac{1}{k^2 \lambda_0^2} \left( 1 + \frac{2}{U} \right) \rightarrow 1$  at  $U \rightarrow \infty$ .

And there is a logarithmic singularity at  $U \rightarrow 1$ .

$\operatorname{Re} E$

