

(a.)  $H = \int \vec{A} \cdot \vec{B} dV$ ,  $\vec{B} = \nabla \times \vec{A}$

H is gauge invariant so  $\frac{d\vec{A}}{dt} = -\vec{E} - \nabla\phi$

$$\frac{dH}{dt} = \int \frac{d\vec{A}}{dt} \cdot \vec{B} + \vec{A} \cdot \frac{d\vec{B}}{dt} dV$$

$$= \int dV \quad -\vec{E} \cdot \vec{B} - \underbrace{\nabla\phi \cdot \vec{B}} - \underbrace{\vec{A} \cdot \nabla \times \vec{E}}$$

$$- \vec{A} \cdot \nabla \times \vec{E} = \nabla \cdot (\vec{A} \times \vec{E}) - \vec{E} \cdot \nabla \times \vec{A} = \nabla \cdot (\vec{A} \times \vec{E}) - \vec{E} \cdot \vec{B}$$

$$-\nabla\phi \cdot \vec{B} = -\nabla \cdot (\phi \vec{B}) + \phi \nabla \cdot \vec{B}$$

$$\frac{dH}{dt} = \int dV \quad -2\vec{E} \cdot \vec{B} + \nabla \cdot (\vec{A} \times \vec{E} - \phi \vec{B})$$

$$= \int dV \quad -2\vec{E} \cdot \vec{B} + \int dS \quad \vec{A} \times \vec{E} - \phi \vec{B}$$

= 0 from Ohm's Law.  
 $\vec{E} = -\vec{v} \times \vec{B}$

= 0, vanishing boundary terms

$\therefore \frac{dH}{dt} = 0 \Rightarrow$  helicity is conserved

$$(b.) SW = \frac{1}{2} \int dV \cdot \left[ \underbrace{\frac{Q_{\perp}^2}{2}}_{(1)} + \frac{B^2}{\mu_0} \underbrace{(\nabla \cdot \vec{\xi}_{\perp} + 2\vec{\xi}_{\perp} \cdot \vec{k})^2}_{(2)} + \underbrace{\gamma P (\nabla \cdot \vec{\xi})^2}_{(3)} - 2 \underbrace{(\vec{\xi}_{\perp} \cdot \nabla P) (\vec{\xi}_{\perp}^* \cdot \vec{k})}_{(4)} - \underbrace{\mu(\vec{r}) (\vec{\xi}_{\perp}^* \times \vec{B}) \cdot Q_{\perp}}_{(5)} \right]$$

(1), (2) and (3) are always positive, thus always stabilizing

(1): bending B    (2): compressing B    (3): compressing plasma

(4) and (5) can have either sign  $\rightarrow$  potentially destabilizing

(4): pressure-driven instability    (5): current-driven instability.

(c.) Ideal force-free equilibria implies  $\vec{E} = -\vec{u} \times \vec{B}$ ,  $\vec{j} \times \vec{B} = 0$   
so we can drop term (4) because  $\nabla p = 0$

the only destabilizing term is  $-\nu(\vec{r}) (\vec{\xi}_{\perp}^* \times \vec{B}) \cdot \vec{Q}_{\perp}$ , so a sufficient condition for stability is

$$\delta W_5 = \frac{1}{2} \int dV -\nu(\vec{r}) (\vec{\xi}_{\perp}^* \times \vec{B}) \cdot \vec{Q}_{\perp} \geq 0$$

$$\vec{Q} = \nabla \times (\vec{\xi}_{\perp} \times \vec{B}) = \nabla \times \vec{A}_{\perp} \quad \text{so} \quad \vec{A}_{\perp} = \vec{\xi}_{\perp} \times \vec{B} \quad \text{so} \quad \vec{A}_{\parallel} = 0$$

$$\text{so } \delta W_5 = \frac{1}{2} \int dV -\nu(\vec{r}) [\vec{A}_{\perp} \cdot \vec{B}_{\perp}]$$

in part (a) it states "no interceptions of the magnetic field" so  $H = \int dV \vec{A} \cdot \vec{B} = 0$ .

If  $\nu(\vec{r}) = \text{const.}$ , ( $\nu \propto J_{\parallel}/B$ ) then we have

$$\delta W_5 = -\frac{J_{\parallel}}{B} \int dV \vec{A}_{\perp} \cdot \vec{B}_{\perp} = -\frac{J_{\parallel}}{B} \int dV \vec{A}_{\perp} \cdot \vec{B}_{\perp} = 0$$

(since the Helicity of the perturbations is also 0)

So a sufficient stability condition is  $\nu(\vec{r}) = \text{const.}$